

Investigation of Interfacial Charge Separation at PbS QDs/(001) TiO₂ Nanosheets Heterojunction Solar Cell

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In the recent years, the heterojunction solar cells based on quantum dots (QDs) have attracted attention due to strong light absorbing characteristics and the size effect on the bandgap tuning. This paper reports on the kinetics of interfacial charge separation of PbS QDs/(001) TiO₂ nanosheets heterojunction solar cells. PbS QDs are deposited using a bifunctional linker molecule on two different TiO₂ films, i.e., TiO₂ nanosheets (with 001 dominant exposed facet) and TiO₂ nanoparticles (with 101 dominant exposed facet). Upon bandgap excitation, electrons are transferred from the PbS QDs conduction band to the lower lying conduction band of TiO₂. Based on the ultrafast pump-probe laser spectroscopy technique, the kinetics of charge separation is scrutinized at the PbS/TiO₂ interface. The interfacial charge separation at PbS/TiO₂ nanosheets films made of (001) dominant exposed facets is five times faster than that on (101) dominant exposed facets TiO₂ nanoparticles. The quantum yields for charge injection are higher for the (001) TiO₂ nanosheets than the (101) TiO₂ nanoparticles due to enhanced interfacial interaction with (001) surface compared to the (101) nanoparticles. The superior interfacial charge separation at PbS/(001) nanosheets respect to PbS/(101) nanoparticles is consistent with the higher photocurrent and enhanced power conversion efficiency in the PbS QDs/(001) TiO₂ heterojunction solar cell. The use of (001) TiO₂ nanosheets can be a better alternative to conventional mesoporous TiO₂ films in QD heterojunction solar cells and perovskites-based heterojunction solar cells.

region. PbS QDs have high dielectric constant, small electron, and hole effective mass and large Bohr radius of 18 nm.^[1–3] The optical absorption by the bandedge excitonic transition can be tuned between 0.3 and 2 eV with a large absorption cross-section (10⁻¹⁵ cm²) and high luminescence quantum efficiency of about 80%.^[4] The tunable bandgap over wide energy range is one great advantage of applying QDs as light-harvesting materials in solar energy conversion systems.^[5] Size-dependent optical properties, high optical extinction coefficient, and new concepts of multiple exciton generation^[6] in QDs or quantum wells to overcome Shockley-Queisser limit for power conversion efficiency are unique advantages of using QDs over other systems for energy-conversion applications.

Several possible structures have been shown for QDs-based solar cell, sensitized QDs solar cells, QD Schottky solar cells, and QDs heterojunction solar cells. The structure of QD heterojunction solar cell composed of mesoscopic TiO₂ bottom layer, which acts as electron collector layer. The top layer is made of QDs that absorb the light. A metal contact usually gold or

silver is deposited on top of the QDs layer to collect the holes. Recent investigations on depleted QDs heterojunction solar cells include varying the size of the QDs,^[7,8] application of TiO₂ or ZnO thin films,^[9–11] surface treatments of the QDs.^[5] Moreover new morphologies of the TiO₂ in QDs-based solar cells have been investigated such as TiO₂ nanosheets with (001) facets,^[12] nanofibers,^[13,14] or nanotubes,^[5] which show superior behavior as electron acceptor layer in contact with dye or QDs as sensitizers. Recently organo-lead halide perovskite was demonstrated to be an efficient light harvester in the solar cell, in one of the pioneer works, the perovskite was used on TiO₂ with (001) dominant facets.^[15] Although the application of dye molecules^[16] and QDs^[12] with (001) TiO₂ nanosheets improved the cell performance, the kinetics of interfacial charge separation at the QD or the dye with (001) TiO₂ nanosheets has never been investigated.

The heterogeneous charge-separation process is a key process, which results in the current generation through the cell. Kamat and co-workers^[17,18] studied the electron injection process from CdSe QDs to TiO₂ nanoparticles. They demonstrated that the vectorial electron-transfer rate is affected by the QDs particle size and the conduction band energy position. Due to the relative smaller electron effective mass than holes effective

1. Introduction

Quantum dots (QDs) have optical bandgaps that are tunable across a wide range of energy levels by changing the QD size. This is in contrast to bulk semiconductor materials, where the bandgap is fixed by the choice of material composition. Semiconductor QDs of the group IV-VI, such as PbS and PbSe are showing strong absorption in the visible and near-infrared

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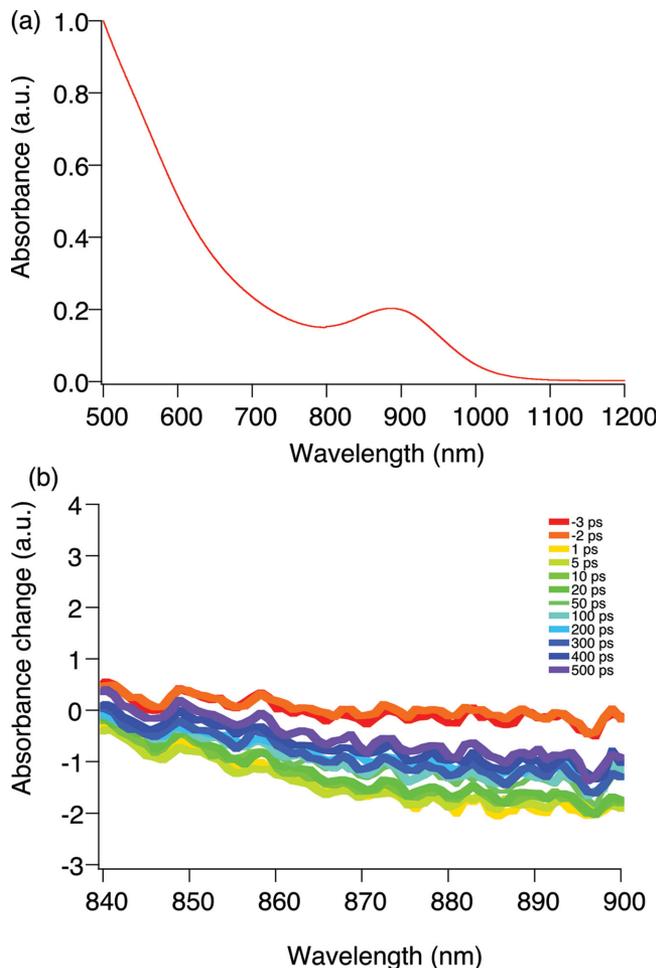


Figure 3. A) Absorption spectra of the PbS nanoparticles used in this study. B) Transient supercontinuum spectrum on PbS/TiO₂ heterojunction film. PbS QDs are excited at 520 nm, and the transient spectrum is shown at wavelengths from 840 to 900 nm centered the excitonic absorption peak. The negative traces show the absorption bleach formation and recovery in QDs due to laser-induced exciton formation and exciton dissociation via electron injection. Traces are recorded from excitation time up to 500 ps after excitation.

conduction band. The conduction band position in QDs is size dependent and is very much affected by quantization. In this study, we used 3 nm PbS QDs (corresponding to band gap of 1.38 eV) linked to TiO₂ particles with different dominant crystallographic facets. **Figure 3A** showed typical absorption spectra of the studied PbS QDs; the first excitonic absorption peak is observed around 880 nm.

To determine the electron injection rate from PbS QDs into the (001) TiO₂ nanosheets and to the (101) TiO₂ NPs, time-resolved transient absorption measurements were performed. The PbS QDs were photo-excited with ultra short 520 nm laser pulses. Upon excitation, electron and hole accumulate in the conduction band and valence band, which induce a decrease of the absorption (bleaching) of QDs at wavelengths corresponding to the exciton absorption band. This absorption feature is in the wavelength range of 840 nm–1 μ m. Transient absorption spectra in the NIR are displayed in **Figure 3B**. Monitoring the

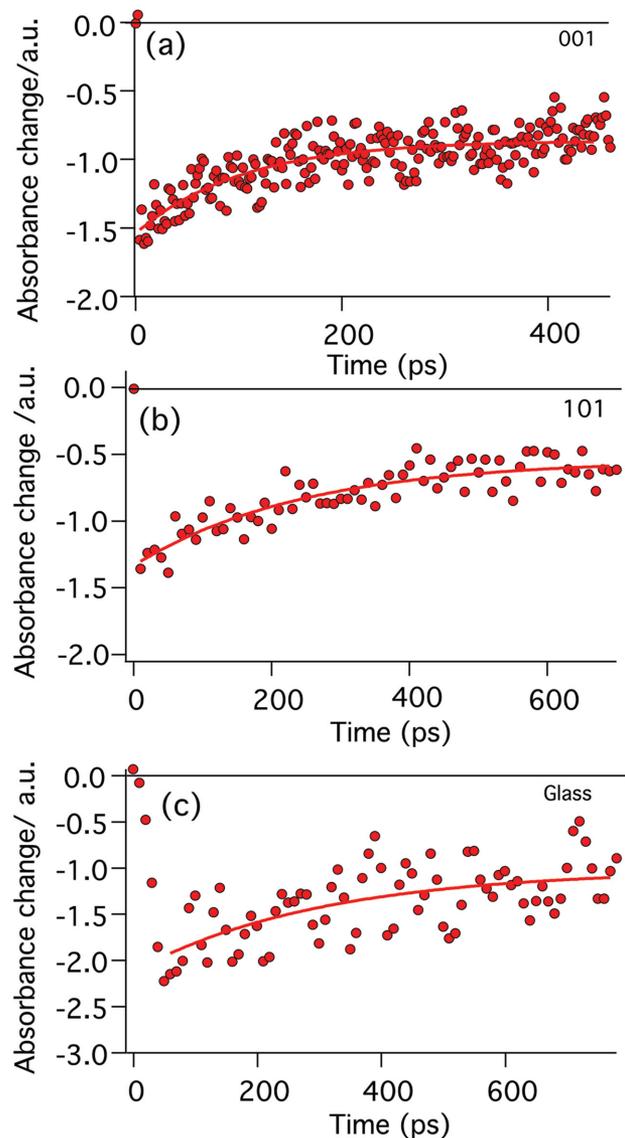


Figure 4. Bleaching dynamics for first excitonic transition (1S_h-1S_e) recorded at 870 nm following 520 nm excitation of PbS QD anchored on TiO₂ (001) nanosheets, (101) nanoparticles films and on top of a glass substrate. The absorbance bleach recovery is fitted by single exponential function. The solid lines represent the single exponential fitting. The fitting parameters are listed in Table 2.

time evolution of the absorption bleaching gives information on charge-separation and recombination processes.

Figure 4 compares the transient absorption at 870 nm of the PbS QDs deposited on TiO₂ nanosheets with (001) dominant facets and TiO₂ NPs with (101) dominant facets. The PbS QDs and sample preparation are identical in both samples. These measurements were compared with that of PbS QDs deposited on glass prepared at identical experimental condition.

The transient absorption decay kinetic for each sample was fitted by single exponential decay function with exponential rate constant (τ^{-1}) and characteristic lifetime (τ) listed in **Table 1**.

For PbS QDs film on a glass substrate, the first-order rate constant (τ^{-1}) is $3.3 \times 10^9 \text{ s}^{-1}$ and increases to $4.1 \times 10^9 \text{ s}^{-1}$ and $9.8 \times 10^9 \text{ s}^{-1}$ when the QDs are anchored on the (101) and (001)

Table 1. Fitting parameters of the transient absorbance signals. τ^{-1} is the first-order rate constant of single exponential fit function to the transient absorptions signals. k_{et} is the interfacial electron-transfer deduced from Equation (1). τ_{et} is the characteristic time constant of interfacial electron-transfer process.

Sample	τ^{-1} [s ⁻¹]	k_{et} [s ⁻¹]	τ_{et} [s]
PbS/TiO ₂ (001)	9.8×10^9	6.5×10^9	154×10^{-12}
PbS/TiO ₂ (101)	4.1×10^9	0.8×10^9	800×10^{-12}
PbS/Glass	3.3×10^9	–	–

TiO₂ films, respectively. For PbS QDs deposited on glass, the recovery of the decay kinetics gives information on charge carrier processes in QDs, such as trapping of photoinduced carriers within the QDs,^[18] followed by radiative recombination kinetic at longer time scale. When PbS QDs are in the vicinity of TiO₂ with lower lying conduction band, time-resolved signal is a convolution of both the interfacial electron-transfer rate and exciton recombination or charge trapping in QDs. Therefore, the acceleration in the observed rate constant is due to the additional charge-transfer pathway from QDs into the TiO₂. Therefore, the charge transfer rate constant can be estimated by Equation (1), with assuming that the entire carrier processes in QDs/glass sample and QDs/TiO₂ films are identical except the electron-transfer process.

$$k_{\text{et}} = \frac{1}{\tau(\text{PbS} + \text{TiO}_2)} - \frac{1}{\tau(\text{PbS})} \quad (1)$$

Using the rate constants values listed in Table 1, we obtained electron-transfer time constants of 800 ps and 154 ps for (101) and (001) TiO₂ films, respectively (see Figure 2B). These results indicate that there is a noticeable difference in electron injection time from PbS QDs to TiO₂ nanostructured films with different morphologies and crystallographic facets.

In order to learn the difference in the photovoltaic performance of these two TiO₂ crystallographic facets, solar cells using the studied PbS QDs were fabricated. The current–voltage characteristic of the cells made with (001) TiO₂ nanosheets and standard cells made of TiO₂ NPs with dominant exposed (101) facets (Dyesol paste DS90) are presented in Figure 5. The photovoltaic parameters of these cells are summarized in Table 1.

We observe that the TiO₂ nanosheets give higher short-circuit photocurrent than the standard TiO₂ NPs and better power conversion efficiency. These results are in agreement with previous studies that reported a better photovoltaic performance in (001) nanosheets than standard (101) nanoparticles, when the films are applied in both dye and QD-sensitized device configurations.^[12,16] Our observation on the better interfacial charge separation at PbS/(001) nanosheets is consistent with the better photocurrent generation in these cells respect to the classical (101) nanoparticles-based devices. Studies by Kavan and co-workers^[16] show that the electron transport is the same in these two film morphologies. Therefore, the faster charge separation kinetic for the nanosheets with (001) dominant facets supports the higher photocurrent achieved in the case of PbS QDs/TiO₂ nanosheets heterojunction solar cells compare to PbS QDs / TiO₂ NPs heterojunction solar cells.

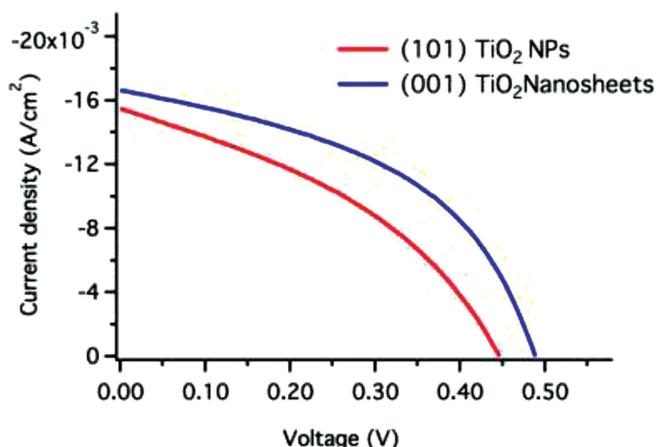


Figure 5. J–V curves of cells made of (001) TiO₂ nanosheets and standard cells made of TiO₂ NPs with dominantly exposed (101) facets.

It should be noted that the PbS QDs are attached to TiO₂ nanoparticles through a bifunctional linker molecule (MPA) that has a thiol group (SH) on one end and carboxylic group on the other side for anchoring to TiO₂. This molecule may contribute to additional deactivation pathway and also decrease the injection rate by increasing the donor–acceptor distance for electron transfer. Moreover, the thiol head groups can act as a hole scavenger if the oxidation potential is appropriate for hole-transfer reaction. The oxidation potential of thiol ligands is located between –5.5 and –6.0 eV versus vacuum.^[19] The HOMO level of the PbS QDs used in this study located at –5.1 eV (see Figure 2B); consequently the hole-transfer reaction is not favorable to occur in this system.

To rationalize the observed faster charge injection in DSCs based on nanosheets respect to nanoparticles, the QD anchoring geometry should be considered. For the electron injection into a continuum of acceptor states implies that the nuclear factor in Fermi golden rule is very large. For a through space tunneling, the electron-transfer rate is controlled by the electronic coupling between the donor and the acceptor states through ($\exp^{-\beta r}$), where r is the distance and β is the damping factor ($\beta \approx 1.2 \text{ \AA}^{-1}$). The observed increase of the electron-transfer rate constant by a factor of 5, corresponds then to a decrease of the electron-transfer distance of about 1.4 Å, which can originate from different association geometry of MPA-capped PbS QDs on two different TiO₂ films. It can be hypothesized that the more closely exposed Ti_{5c} sites at (001) crystalline surface provide a better anchoring of the QDs to the (001) films via MPA linker molecule. A more precise theoretical modeling of the association geometry of QDs to the TiO₂ (001) and (101) surface via MPA linker molecule would be the subject of a separate study. In our previous publication on DSCs-based nanosheets, it was observed that the back electron transfer was by a factor of 6 slower for the TiO₂ (001) nanosheets compared to the same process on TiO₂ (101) NPs, that was rationalized by different association of carboxylic anchoring group on the two films of TiO₂. The different anchoring geometry of carboxylic on the two crystalline films was also proved by theoretical investigations using DFT techniques by De Angelis and co-workers.^[24]

Table 2. Photovoltaic parameters for the (001) TiO₂ nanosheets cells and standard cells made of TiO₂ NPs with dominantly exposed (101) facets.

TiO ₂ Sample	QDs E _g [eV]	V _{oc} [V]	J _{sc} [mA cm ⁻²]	FF	η [%]	BET [m ² g ⁻¹]
TiO ₂ 18 nm NPs (DS90)	1.38	0.44	15.48	0.38	3.32	75
TiO ₂ nanosheets (001) dominant facets	1.38	0.49	16.63	0.46	4.17	122

In addition to photocurrent enhancement, the open circuit voltage was increased by 50 mV for the PbS/(001) TiO₂ solar cell compare to the PbS/(101) TiO₂ solar cell (see **Table 2**). This is again consistent with our previous work^[16] on Ru-based DSCs where higher Voc was also achieved for Ru-sensitized (001) TiO₂ nanosheet photoanode respect to Ru-sensitized (101) TiO₂ NPs. The higher Voc was rationalized by a negative shift in the conduction band of the TiO₂ nanosheets with (001) dominant facets compared to the TiO₂ NPs with (101) dominant facets, which was later proved by the theoretical analysis based on density functional theory calculations by De Angelis and co-workers.^[24]

Finally, the PbS/(001) TiO₂ heterojunction solar cell demonstrates better quantum yield (Φ_{inj}) for charge injection. The quantum yield for charge injection into TiO₂ depends upon the kinetic competition between interfacial electron-transfer rate (*k_{et}*) and recombination rate (*k_r*) within the QDs given by Equation (2). The recombination rate constant, *k_r* ≈ τ⁻¹, was measured for the PbS QDs on glass substrate. Therefore, the quantum yield of charge injection was calculated to be 0.75 for (001) nanosheets and 0.55 for (101) NPs, which agree well with the better PV performance of the PbS/(001) TiO₂ heterojunction solar cells.

$$\Phi_{\text{inj}} = \frac{k_{\text{et}}}{k_{\text{r}} + k_{\text{et}}} \quad (2)$$

To summarize, this report provides complementary understanding to the existing studies on charge separation in QDs based solar cells, it uses pump probe laser spectroscopy to elucidate the interfacial charge separation at different TiO₂ crystallographic facets in heterojunction PbS QDs solar cell.

3. Conclusion

This work describes for the first time the rate of electron transfer from PbS QDs to different TiO₂ crystallographic facets in PbS QDs heterojunction solar cell. Two different TiO₂ dominant facets were studied, the TiO₂ (101) exposed facets and the TiO₂ (001) exposed facets. The electron-transfer rate from the PbS QDs to the (001) exposed facets was faster by a factor of 5 than the electron-transfer rate to the (101) exposed facets. This observation is consistent with the higher current density and the better power conversion efficiency observed in the PbS QDs/(001) TiO₂ heterojunction solar cells. Moreover, the quantum yield for charge injection was higher for the (001) dominant facets compare to the (101) dominant facets, which further support the better PV performance. This work demonstrates that the TiO₂ nanosheets with exposed (001) facets can

be a better alternative to the conventional (101) TiO₂ NPs in QDs and possibly in perovskite-based solar cells.

4. Experimental Section

Synthesis and Purification of TiO₂ Nanosheets: The synthesis of TiO₂ nanosheets followed a typical experimental procedure.^[15] 10 mL of Ti(OBu)₄ (98%) and 0.8 mL of hydrofluoric acid (47%) solution were mixed in 150 mL dried Teflon autoclave, which was kept at 180 °C for 24 h to yield the nanosheets. After the reaction mixture was cooled to room temperature, the white powder was separated by high-speed centrifugation and washed with ethanol followed by distilled water for several times.

Caution: Hydrofluoric acid is extremely corrosive and a contact poison, it should be handled with extreme care! Hydrofluoric acid solution is stored in Teflon containers in use.

Methods and Device Fabrication: Colloidal PbS QDs were purchased from Evident Technology and were kept under inert atmosphere. During device fabrication, a thin blocking layer of compact TiO₂ was first deposited by spray pyrolysis onto a precleaned FTO glass substrate using a solution of titanium diisopropoxide bis(acetylacetonate) in ethanol as precursor. Finally, a layer of TiO₂ nanosheets was deposited by spin-coating (4000 rpm for 10 s) from dilute aqueous solution containing a mixture of TiO₂ nanosheets and ethanol (volume ratio of 1:3). This results in a nanosheet titania film with 300 ± 50 nm thickness. The procedure for the deposition of the PbS QDs film is adapted from the one reported by Sargent and co-workers.^[23] The PbS QDs were subsequently deposited layer by layer on the porous TiO₂ film by spin coating a 50 mg mL⁻¹ solution in octane. Each layer was casted at a spinning rate of 2500 rpm applied for 10 s and treated thereafter briefly with a solution of 10% 3-mercaptopropionic acid (MPA, ≥99.0% from Sigma-Aldrich) in methanol using again 2500 rpm rotational speed for 10 s. This treatment displaced the oleate ligand and rendered the QD insoluble, which allowed thin films of 300 nm thicknesses to be created using 12 successive deposition cycles. Each layer was rinsed with anhydrous methanol (purchased from Sigma-Aldrich) to remove excess of MPA and PbS QDs. Finally, gold back contact of ≈100 nm thicknesses was deposited by evaporation through a shadow mask. The device was then completed by encapsulation in an argon atmosphere. The device was encapsulated by gluing a 2-mm thick glass plate on top of the active area of the device using a frame of hot melt Surlyn (25 μm thick) as a sealant and spacer. After encapsulation, the device was stable in ambient air.

Photovoltaic Characterization: Photovoltaic measurements employed an AM 1.5 solar simulator equipped with a 450 W xenon lamp (Model No. 81172, Oriel). Its power output was adjusted to match AM 1.5 global sunlight (100 mW cm⁻²) by using a reference Si photodiode equipped with an IR-cutoff filter (KG-3, Schott) in order to reduce the mismatch between the simulated light and AM 1.5 (in the region of 350–750 nm) to less than 2% with measurements verified at two PV calibration laboratories [ISE (Germany), NREL (USA)]. I–V curves were obtained by applying an external bias to the cell and measuring the generated photocurrent with a Keithley model 2400 digital source meter. The voltage step and delay time of photocurrent were 10 mV and 40 ms, respectively. Photovoltaic performance was measured by using a metal mask with an aperture area of 0.12 cm².

Pump-Probe Laser Spectroscopy: Time-resolved transient absorption measurements were performed on PbS QDs (0.5 μm thick)/TiO₂

mesoporous heterojunction films. The pump-probe spectroscopy technique used a CPA-2001, 1 kHz repetition rate, Ti: Sapphire amplified femtosecond laser (Clark-MXR), delivering 1 mJ pulses of 150 fs width at a central wavelength of 775 nm. The excitation beam was provided by pumping a double-stage noncollinear optical parametric amplifier (NOPA). The NOPA was tuned at 530 nm to generate pulses of approximately 8 μ J energy, which were then compressed in an SF10-glass prism pair down to a duration of less than 60 fs (fwhm). As the probe beam, a white light continuum was generated by focusing a small fraction of 775 nm pulses from the CPA laser system by a $f = 50$ mm lens into a 2-mm-thick CaF₂ rotating disc. The time delay between excitation and probe pulses was controlled precisely by means of a motorized translation stage (M531 PD, Physik Instrument) equipped with a retro-reflector (49-668, Edmund Optics). At the sample, pulse energy was decreased down to a few microjoules for the pump and to less than 1 μ J for the probe beam. Transient absorption spectra were calculated from shot-by-shot simultaneous measurements of signal and reference beams using a pair of spectrographs (Princeton Instruments, SpectraPro 2150) equipped with 512 \times 58 pixels back-thinned CCD detectors (Hamamatsu S7030-0906) and by chopping the probe beam at a frequency of 500 Hz. Time-resolved experiments were applied on sealed samples.

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